## First Measurement of

$$
\operatorname{BR}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{c J}\right)
$$

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## $c \bar{c}$ Spectrum



| Ig(Jpc) | Name | Mass | Width |
| :--- | :--- | :--- | :--- |
|  |  | MeV | MeV |
| $0-(1--)$ | Psi" $^{\prime \prime}$ | 3773 | 25.3 |
| $0-(1--)$ | Psi' $^{\prime}$ | 3686 | 0.28 |
| $0+(2++)$ | $\mathrm{Xc2}$ | 3556 | 2.00 |
| $0+(1++)$ | Xc1 | 3510 | 0.88 |
| $0+(0++)$ | Xc0 | 3415 | 14.90 |
| $0-(1--)$ | $\mathrm{J} / \mathrm{psi}$ | 3096 | 0.09 |
| $0+(0-+)$ | EtaC | 2979 | 17.30 |

## Motivation

1. Previous measurements did not produce significant signals
2. Impact on interpretation of $X(3872)$
3. Validation of Potential Model calculations above the open flavor threshold

## Lack of significant previous measurements

- Before this result, no significant measurement of $\mathrm{BR}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{c J}\right)$
(Unpublished results )
- ~3/pb by MARKII
- ~2/pb by Crystal-Ball
- ~9/pb by MARKIII
- 281/pb; ~30times larger sample from CLEO-c


## What do we know about $X(3872)$

## A.Observed in X(3872) $\rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi$ decay

B.Mass (3871.9 $\pm 0.5$ ) MeV
-just below the DD* threshold
C. Width < 2.3 MeV
-Surprisingly small, since the mass is well above the $D \bar{D}$ threshold
D. No radiative transitions to $\chi_{c J}$ states have been observed
-just an upper limit for decay to $\gamma \chi_{c 1,2}$
$\frac{\Gamma\left(\mathrm{X}(3872) \rightarrow \gamma \chi_{\mathrm{cl}}\right)}{\Gamma\left(\mathrm{X}(3872) \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)}<0.9$ Belle

$$
\frac{\Gamma\left(\mathrm{X}(3872) \rightarrow \gamma \chi_{\mathrm{c} 2}\right)}{\Gamma\left(\mathrm{X}(3872) \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)}<1.1 \text { Belle }
$$

## Possible interpretations of $X(3872)$



Existence of hybrids and bound states of mesons has not been experimentally proven so far

## How our measurement helps?

- Conventional charmonium candidates for $X(3872)$
- $\psi_{2}\left(1^{3} D_{2}\right), h_{c}^{\prime}\left(2^{1} P_{1}\right), \psi_{3}\left(1^{3} D_{3}\right)$
( $C=-1$ )
- $\eta_{c 2}\left(1^{1} D_{2}\right), \chi_{c 1}{ }^{\prime}\left(2^{3} P_{1}\right), \eta_{c}{ }^{\prime \prime}\left(3^{1} S_{0}\right)$,
( $C=+1$ )
- Nonrelativistic-case:<1D|r|1P> is independent of $J$.
- $\Gamma_{J}=4 /{ }_{3} e^{2} \alpha E_{\gamma}{ }^{3} C_{J}\left|<1^{3} D\right| r\left|1^{3} P>\right|^{2}$
- $1^{3} D_{J} \rightarrow \gamma \chi_{c 1}$ can be measured for different $J$ provided one is known
- Measuring $1^{3} D_{1}\left(\psi^{\prime \prime}\right) \rightarrow \gamma \chi_{c 1}$ can shed some light on $1^{3} D_{2,3}\left(\psi_{2}, \psi_{3}\right) \rightarrow \gamma \chi_{c 1}$


## Validation of Potential Model for $\psi^{\prime \prime}$

- Is $\psi^{\prime \prime}$ a pure cc state?
- Strong indications that $X(3872)$ is not
- May be all states above the flavor thresholds have complex nature?
- Radiative transitions are a good probe
- Pure cc state: mostly $1^{3} D_{1}$ (small contribution from $2^{3} S_{1}$ )
- J-dependence of $\Gamma\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{c J}\right)$ is well predicted
- Are relativistic corrections important?
- Are coupled channel effects
$\bar{C} \rightarrow D \bar{D} \rightarrow c \bar{C}$ important?


## CLEO-c

$N_{\psi^{\prime}}=1.5 \mathrm{M}$<br>$\mathrm{N}_{\psi^{\prime \prime}}=1.8 \mathrm{M}$



## Analysis

- Method1
- $\psi(3770) \rightarrow \gamma \chi_{C J} \rightarrow \gamma J / \psi$ $\rightarrow \gamma \gamma \mathrm{l}^{+} \mathrm{l}^{-}$
- Select events with exactly 2 photons and 2 leptons with no net charge:
- No other photon with E> 60 MeV
- $\left|P_{\text {tot }}\right|<50 \mathrm{MeV}$
- $\left|E_{\text {tot }}-E_{\| \mid}-E_{J / \psi}\right|<40 \mathrm{MeV}$
- Electron
- $E / p>0.7$
- Muon
- $0.15<E<0.55 \mathrm{GeV}$
- Signal variable: Energy of lower energy photon
- Method2
- $\psi(3770) \rightarrow \gamma \chi_{c J} \rightarrow \gamma(2 K, 2 K 2 \pi$, $4 \pi, 6 \pi)$
- Select events with exactly 2,4,6 charged hadrons and a photon:
- Highest energy neutral cluster in the calorimeter is the photon candidate
$-\left|P_{\text {tot }} i-P_{c m} i\right|<30 \mathrm{MeV}$
$i=x, y, z-E_{c m} \mid<30 \mathrm{MeV}$
- Kaon
- Combined log-likelihood > 0
- $\left|\sigma_{\mathrm{K}}\right|<3$
- Pion
- Not a kaon
- $\left|\sigma_{\pi}\right|<3$
- Signal variable: Photon energy


## Kinematic fitting

1. Constrain total energy and momentum to the expected values.
2. For $\left.\gamma \gamma l^{+}\right|^{-}$also constrain mass of $\left.I^{+}\right|^{-}$to the $\mathrm{J} / \psi$ mass.

Demonstration on $\psi^{\prime}$ data



## $\psi(2 S)$ background in $\psi(3770)$ data

- ISR production of $\psi(2 S)$ at $E_{c m}=3770 \mathrm{MeV}$
- $\mathrm{e}+\mathrm{e}-\rightarrow \gamma \psi(2 \mathrm{~S})$

$$
\begin{aligned}
& \text { • } \psi(2 S) \rightarrow \gamma \chi_{c J} \rightarrow \gamma J / \psi \rightarrow \gamma \gamma I^{+I-} \\
& \cdot \\
& \cdot \\
& \hline(2 S) \rightarrow \gamma \chi_{c J} \rightarrow \gamma(2 K, 2 K 2 \pi 2 K, 4 \pi, 6 \pi)
\end{aligned}
$$

- $E_{\gamma}^{\text {ISR }} \sim 84 \mathrm{MeV}$ for $\psi(2 S)$ produced with its nominal mass:
- Selection criteria and kinematic fitting gets rid of this background ( $E_{\gamma}^{\text {ISR }}$ forced to be less than about 40 MeV )
- Radiative flux peaks for $E_{\gamma}^{\text {ISR }} \rightarrow 0$ making the remaining background indistinguishable from the signal:
- Estimate this background using $\psi(2 S)$ measurements and theoretical formulae extrapolating the rate to the ISR peak


## $\left.\left.\psi^{\prime \prime} \rightarrow \gamma \gamma\right|^{+}\right|^{-}$

Separate $\mu \mu$ and ee data because of very different background level but fit them simultaneously

Number of events for $\psi^{\prime \prime}$ $A 0=22 \pm 9$
A1 $=53 \pm 10$
$A 2=0 \pm 2.9$
RR from $\psi^{\prime}$
AO $=11.7$
A1 $=20.0$
$A 2=0.6$


## Cross check of $\left.\left.\psi^{\prime \prime} \rightarrow \gamma \gamma\right|^{+}\right|^{-}$ analysis by $\left.\left.\psi^{\prime} \rightarrow \gamma \gamma\right|^{+}\right|^{-}$



Fit to gaussian shapes with linear background
Number of events for $\psi$ '
A1 $=1718 \pm 42$
A2 $=835 \pm 30$

|  | Our measurements of <br> $\operatorname{BR}\left(\psi^{\prime} \rightarrow \gamma c_{c J}\right)(\%)$ | previous CLEO-c <br> measurement using a <br> different technique |
| :---: | :---: | :---: |
| $\mathrm{J}=2$ | $1.84 \pm 0.07$ | $1.81 \pm 0.06$ |
| $\mathrm{~J}=1$ | $3.53 \pm 0.09$ | $3.50 \pm 0.08$ |

Ey low (MeV)

## Results for $\gamma \gamma l^{+\mid-}$Analysis

$$
\operatorname{BR}\left(\psi^{\prime \prime} \rightarrow \gamma x_{c J} \rightarrow \gamma l^{+} l^{-}\right)=\frac{N_{\text {events }}\left(\psi^{\prime \prime} \rightarrow m^{+}+\right)}{\varepsilon_{\psi^{\prime \prime} \rightarrow m^{+t}} \times N_{\varphi^{\prime \prime}}}
$$

|  | Results for $\psi^{\prime \prime} \rightarrow \gamma \chi_{c J}$ |  |  |
| :---: | :---: | :---: | :---: |
|  | $J=2$ | $J=1$ | $J=0$ |
| $\varepsilon(\%)$ | 18 | 23 | 20 |
| Branching Ratio: <br> BR (10-3) | $<0.9$ | $2.8 \pm 0.5 \pm 0.4$ | $<44$ |

- $\mathrm{B}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{c 0}\right)$ is predicted to be the largest, but the small $B\left(\chi_{c 0} \rightarrow \gamma J / \psi\right)$ limited our sensitivity
- In order to measure we $\mathrm{B}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{c 0}\right)$ turned to hadronic decays of $\chi_{c J}$.


## Technique for $2^{\text {nd }}$ Method

$$
\begin{aligned}
& \mathrm{R}=\frac{\mathrm{BR}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{\mathrm{cJ}}, \chi_{\mathrm{cI}} \rightarrow \text { final state }\right)}{\operatorname{BR}\left(\psi^{\prime} \rightarrow \gamma \chi_{\mathrm{CI}}, \chi_{\mathrm{cI}} \rightarrow \text { final state }\right)} \\
& =\frac{N_{\text {events }}\left(\psi^{\prime \prime} \rightarrow \text { final state }\right)}{N_{\text {events }}^{\left(\psi^{\prime} \rightarrow \text { final state }\right)}} \times \frac{\varepsilon_{\left(\psi^{\prime} \rightarrow \text { final state }\right)}}{\varepsilon_{\left(\psi^{\prime \prime} \rightarrow \text { final state }\right)}} \times \frac{N^{\prime}}{N^{\prime \prime}} \\
& \operatorname{BR}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{C J}\right)=\mathrm{R} \times \operatorname{BR}\left(\psi^{\prime} \rightarrow \gamma \chi_{c J}\right) \\
& \text { PR D70, 112002(2004) }
\end{aligned}
$$

## $\psi^{\prime \prime} \rightarrow \gamma 4 \pi($ left $), \gamma 2 \mathrm{~K} 2 \pi$ (right)



## $\psi^{\prime \prime} \rightarrow \gamma 6 \pi$ (left), $\gamma 2 \mathrm{~K}$ (right)



## Combined plots for four hadronic decay modes



Number of events for $\psi^{\prime}$
AO $=2816 \pm 58$
A1 $=886 \pm 32$
A2 $=1329 \pm 40$
Sum of fits (3 CBL with a linear background) to individual decay modes

Number of events for $\psi^{\prime \prime}$ AO $=274 \pm 27$ A1 $=54 \pm 17$
$A 2=20 \pm 18$
RR from $\psi^{\prime}$
$A 0=25.2$
A1 $=12.0$
A2=24.9
Sum of fits (6 CBL with quadratic background)

## Results for $\psi^{\prime \prime} \rightarrow \gamma \chi_{c J}$

|  | Results for <br> $\left(10^{-3}\right)$ |  |  |
| :---: | :---: | :---: | :---: |
|  | $\mathrm{BR}\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{C J}\right)$ |  |  |
| 2nd method | $<2.0$ | $3.9 \pm 1.4 \pm 0.6$ | $7.3 \pm 0.7 \pm 0.6$ |
| $\gamma^{\text {nd }}$ | $<0.9$ | $2.8 \pm 0.5 \pm 0.4$ | $<44$ |
| Combined | $<0.9$ | $2.9 \pm 0.5 \pm 0.6$ | $7.3 \pm 0.7 \pm 0.6$ |

Observe significant signal for $\psi^{\prime \prime} \rightarrow \gamma \chi_{c 0,1}$ and set a $90 \%$ C.L. upper limit for $\gamma \chi_{c 2}$.

## Interpretation of $X(3872)$

$$
\begin{aligned}
& \frac{\Gamma\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{\mathrm{cl}}\right)}{\Gamma\left(\psi^{\prime \prime} \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)}=1.56 \pm 0.37 \pm 0.37 \\
& \text { PRL } 96082004 \text { (2006) } \\
& \frac{\Gamma\left(\psi_{2} \rightarrow \gamma \chi_{\mathrm{cl}}\right)}{\Gamma\left(\psi_{2} \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)} \approx(2-3.5) \times \frac{\Gamma\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{\mathrm{cl} 1}\right)}{\Gamma\left(\psi^{\prime \prime} \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)}>2 \times 1.8 \\
& \frac{\Gamma\left(\mathrm{X}(3872) \rightarrow \gamma \chi_{\mathrm{cl}}\right)}{\Gamma\left(\mathrm{X}(3872) \rightarrow \pi^{+} \pi^{-} \mathrm{J} / \psi\right)}<0.9 \text { Belle }
\end{aligned}
$$

## $X(3872)$ is not $1^{3} D_{2}$ !

## Nature of $\psi(3770)$

- Theoretically
- $\Gamma_{J}=4 /\left.{ }_{3} e^{2} \alpha E_{\gamma}{ }^{3} C_{J}\left|<1^{3} D_{1}\right| r\left|1^{3} P_{J}\right\rangle\right|^{2}$
- Non-relativistically $<1^{3} D_{1}|r| 1^{3} P_{J}>$ is $J$ independent
- we can cancel it by calculating the ratios of widths
- J-dependence:
- $C_{\mathrm{J}}=2 / 9,1 / 6$ and $1 / 90$ for $1^{3} \mathrm{D}_{1} \rightarrow 1^{3} \mathrm{P}_{\mathrm{J}} \mathrm{J}=0,1$ and 2
- Measured $E_{\gamma}$
- Thus in non-relativistic limit expect:
- $\Gamma_{0} / \Gamma_{1}=3.2 \quad$ and $\Gamma_{0} / \Gamma_{2} \sim 85$
- Measured:
- $\Gamma_{0} / \Gamma_{1}=2.5 \pm 0.6$ and $\Gamma_{0} / \Gamma_{2}>8$


## Evidence that $\psi^{\prime \prime}$ is predominantly $1^{3} D_{1}$ state

## Beyond the naïve theory

|  | $\Gamma\left(\psi^{\prime \prime} \rightarrow \gamma \chi_{\mathrm{cJ}}\right) \quad(\mathrm{keV})$ |  |  |
| ---: | :---: | :---: | :---: |
|  | $\mathrm{J}=2$ | $\mathrm{~J}=1$ | $\mathrm{~J}=0$ |
| CLEO-c data | $<20$ | $70 \pm 17$ | $172 \pm 30$ |
| Rosner (non-relativistic) | $24 \pm 4$ | $73 \pm 9$ | $523 \pm 12$ |
| Ding-Qin-Chao |  |  |  |
| Non-relativistic | 3.6 | 95 | 312 |
| Relativistic | 3.0 | 72 | 199 |
| Eichten-Lane-Quig9 |  |  |  |
| Non-relativistic | 3.2 | 183 | 254 |
| Coupled-channel corrections | 3.9 | 59 | 225 |
| Barnes-Godfrey-Swanson |  |  |  |
| Non-relativistic | 4.9 | 125 | 403 |
| Relativistic | 3.3 | 77 | 213 |

- Relativistic/coupled-channel corrections in potential model calculations are important for agreement with the data


## Decay width for $\psi(2 S)$

|  | $\Gamma\left(\psi^{\prime} \rightarrow \gamma \chi_{c J}\right)(\mathrm{keV})$ |  |  |
| ---: | :---: | :---: | :---: |
|  | $\mathrm{J}=2$ | $\mathrm{~J}=1$ | $\mathrm{~J}=0$ |
| CLEO data | $27 \pm 4$ | $27 \pm 3$ | $27 \pm 3$ |
| Rosner (non-relativistic) | $35 \pm 1$ | $75 \pm 3$ | $26 \pm 6$ |
| Ding-Qin-Chao |  |  |  |
| Non-relativistic | 42 | 36 | 25 |
| Relativistic | 25 | 28 | 22 |
| Eichten-Lane-Quig9 |  |  | 36 |
| Non-relativistic | 23 | 33 | 38 |
| Coupled-channel corrections | 23 | 32 | 63 |
| Barnes-Godfrey-Swanson |  | 54 | 26 |
| Non-relativistic | 38 | 29 |  |
| Relativistic | 24 |  |  |

- Corrections needed in potential model calculations for agreement with the data of $\psi(2 S)$ as well


## Conclusions

- We have observed $\psi^{\prime \prime} \rightarrow \gamma \chi_{c 0,1}$ for the first time:
- $\quad$ r $l l$ results published in PRL 96, 182002 (2006)
- results for hadronic in PR D, Rapid Communications (hep-ex/0605070)
- In view of our results the $1^{3} D_{2}$ interpretation of $X(3872)$ can be ruled out
- Spin dependence of the observed rates confirms that $\psi(3770)$ is predominantly $1^{3} D_{1}$ state
- Relativistic or couple channel effects are needed for quantitative agreement between potential model calculations and the data


## BACK UP

## Some detector plots



## K-f Effect

- 0-ggll
- DATA: 128.855 .629
- w/k-fit: 128.935 .337
- w/o: 127.577 .082
- 1-4pi
- DATA: 127.575 .245 fixed to MC
- w/ k-fit w/ k-fit: 126.944 .9691 .224
- w/o: 127.116 .1860 .945
- 2-2k2pi
- DATA: 128.455 .422 fixed to MC
- w/k-fit w/ k-fit: 126.834 .8351 .265
- w/o: 126.956 .3391 .009
- 3-6pi
- DATA: 127.945 .172 fixed to MC
- w/k-fit w/ k-fit: 126.914 .9661 .201
- w/o: 127.136 .2630 .932
- 4-2k
- DATA: 127.544 .851 fixed to MC
- w/ k-fit w/ k-fit: 126.874 .5701 .220
- w/o: 126.936 .3050 .986


## 2S1-1D1 mixing

- The measured rate for
 $\mathrm{J}=0$ is much larger than for $J=1$ (which in turn is larger than $\mathrm{J}=2$ ).
- Confirming naïve prediction $B R_{0} \quad B R_{1} \quad B R_{2}$
- Confirming D state

Insensitive to mixing

- Mixing needed to explain large cross-section of $\psi(3770)$ in $e^{+} e^{-}$ experiment
- Effects of mixing on the rates are small
- Can be explored more with better measurement of $J=2$


## General

- mDD*=3871.2 MeV (neutral), 3879.3 MeV (charged)
- mDD $=3729.0 \mathrm{MeV}$ (neutral), 3738.0 MeV (charged)
- Eqn's of k-fit
- $\mathrm{Pcm}=\mathrm{Pl}++\mathrm{Pl}-+\mathrm{Pg} 1+\mathrm{pg} 2$
$-\mathrm{Pcm}=\mathrm{Ph}++\mathrm{Ph}-+\mathrm{Pg}$
- ISR background

$$
\mathrm{N}_{\text {events }}\left(\psi^{\prime} \text { in } \psi^{\prime \prime} \text { from ISR }\right)=B_{\left(\psi^{\prime} \rightarrow \text { final state }\right)} \times \varepsilon_{\left(\psi^{\prime \prime} \rightarrow \text { final state }\right)} \times L_{\psi^{\prime \prime}} \times \Gamma_{e e}\left(\psi^{\prime}\right) \times I(s)
$$

$=\frac{N_{\text {events }}{ }^{\left(\psi^{\prime} \rightarrow \text { final state }\right)}}{\mathrm{N}_{\psi^{\prime}} \times \varepsilon_{\left(\psi^{\prime} \rightarrow \rightarrow \text { final state }\right)}} \times \varepsilon_{\left(\psi^{\prime \prime} \rightarrow \text { final state }\right)} \times L_{\psi^{\prime \prime}} \times \Gamma_{e e}\left(\psi^{\prime}\right) \times I(s)$
$I(s)=\int_{0}^{x} W(s, x), b\left(s^{\prime}(x)\right) F_{X}\left(s^{\prime}(x)\right) d x$
$F\left(s^{\prime}\right)=\left(E_{\gamma}^{I S R}\left(s^{\prime}\right) / E_{\gamma}^{I S R}\left(M_{R}{ }^{2}\right)\right)^{3}$

