



Lecture # 2

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The interactions of photons

- Behavior of photons (**x-rays, γ -rays**) \longrightarrow different from charged particles
- **Photo electric effect**
- **Compton scattering**
- **Pair production**
- **x-rays and γ -rays** are many times more penetrating
- Much **smaller cross-section** relative to electron inelastic collisions
- Beam of photons is **not degraded** \longrightarrow in energy but **attenuated** \longrightarrow in intensity

- Three process remove photons completely from beam → absorption or scattering
- Photon pass straight → retain orig. energy
- However total no. of photons → reduced
- Attenuation → $I(x) = I_0 \exp(-\mu x)$

Incident beam intensity

Absorption coefficient

Absorber thickness

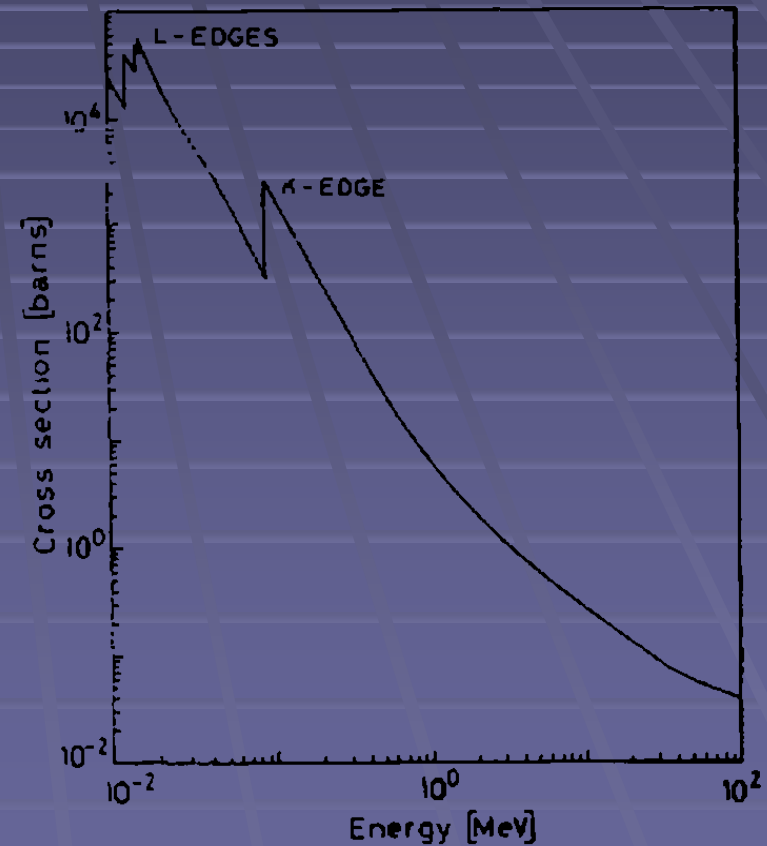
- Absorption coefficient → characteristic of absorbing material → related to total interaction cross-section

Photoelectric effect

- Photoelectric effect → absorption of photon by atomic electron → subsequent ejection of electron from atom
- Energy of electron → $E = h\nu - B.E$
- Photoelectric effect → always occur → on bound electrons
- Nucleus → absorb recoil momentum
- Energies above → highest electron B.E → relatively small
- Increases → as **k-shell** energy approached

- Drops → k-shell electron not available for P.E
- **K-absorption edge**
- Below this energy cross-section rises →
- **L-absorption edge, M-absorption edge**
- Energies above k-shell always **k-electrons** involved in P.E →
- Dependence on **Z**

Calculated photoelectric cross-section for lead



$$\Phi_{photo} = 4\alpha^4 \sqrt{2} Z^5 \Phi_0 (m_e c^2 / h\nu)^{7/2}$$

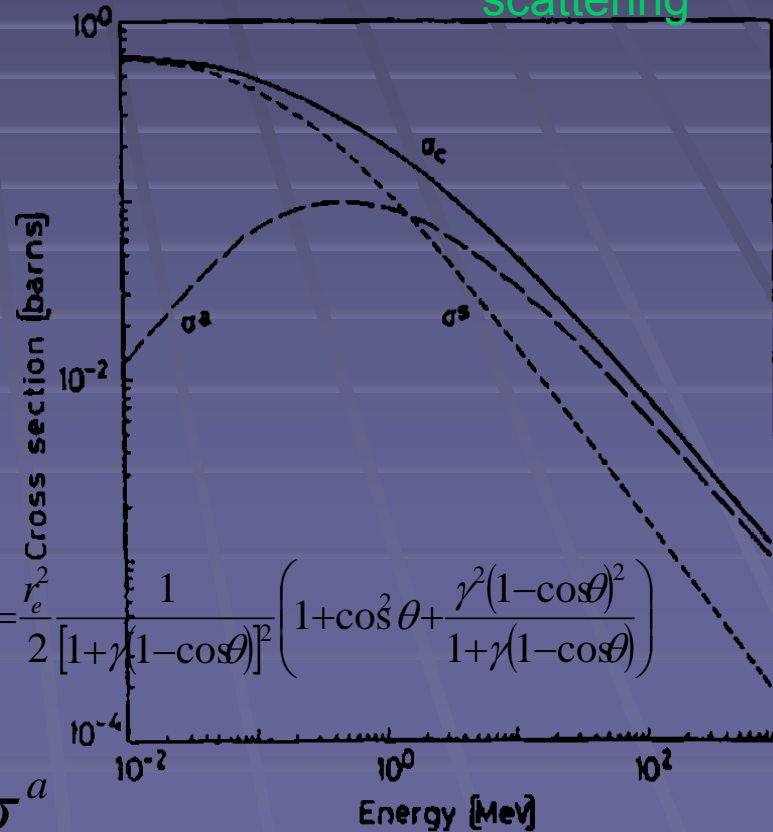
Cross-section by using Born approximation

Compton scattering

- Best understood → process in photon interaction
- Scattering of → photons on free electrons
- Klein-Nishina formula
- **Compton scattered cross-section** → average fraction of total energy contained in scattered photon
- **Compton absorption cross-section** → average energy transferred to recoil electron



Kinematics of Compton scattering



Average energy absorbed by material

$$\sigma_c = \sigma^s + \sigma^a$$

Applying energy and momentum conservation, following relations are obtained

$$h\nu^- = \frac{h\nu}{1 + \gamma(1 - \cos \theta)}$$

$$T = h\nu - h\nu^- = h\nu \frac{\gamma(1 - \cos \theta)}{1 + \gamma(1 - \cos \theta)}$$

$$\cos \theta = 1 - \frac{2}{(1 + \gamma)^2 \tan^2 \phi + 1}$$

Thomson sca. → photon → free electron → classical

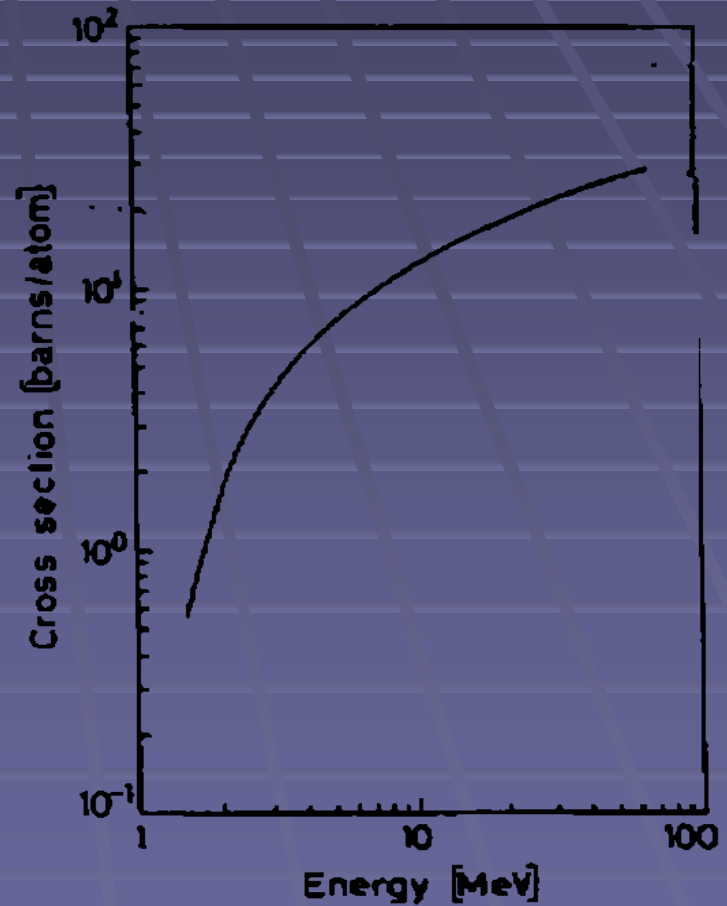
Rayleigh sca. → Photon by → atom as a whole

In both process → no energy transfer → to medium

No excitation → no ionization → only photon direction change

Pair production

- Transformation of photon \rightarrow electron-positron pair
- To conserve momentum \rightarrow in presence of nucleus
- Photon energy \rightarrow 1.022 MeV
- Theoretically pair production \rightarrow related to bremsstrahlung



Pair production cross-section in lead

Introduction

- General characteristics **common** to **detectors**
- Transfer of **part** or **all** of radiation energy to detector mass
- **Charged** particles produce **direct** ionization
- **Neutral** particles produce **indirect** ionization
- Form of converted energy depends on detector and its design
- **Gaseous** detectors are designed to **directly collect** ionization
- In **scintillators** both **ionization** and **excitation** contribute to induce molecular excitations, which results in emission of light

- In photographic emulsion, ionization induces chemical reactions
- Detectors are **electrical** in nature
- Informations are transformed into **electrical pulses**
- Recordable information
- Specific **data formats**



General Characteristics of detectors

- Sensitivity
- Detector response
- Energy resolution: The Fano Factor
- The response function
- Response time
- Detector efficiency
- Dead time

Sensitivity

- The capability to produce a usable signal
 - No detector can be sensitive to all radiations
1. The **cross-section** for ionizing reactions in the detector
 2. The detector **mass**
 3. The inherent detector **noise**
 4. The **protective material** surrounding the sensitive volume

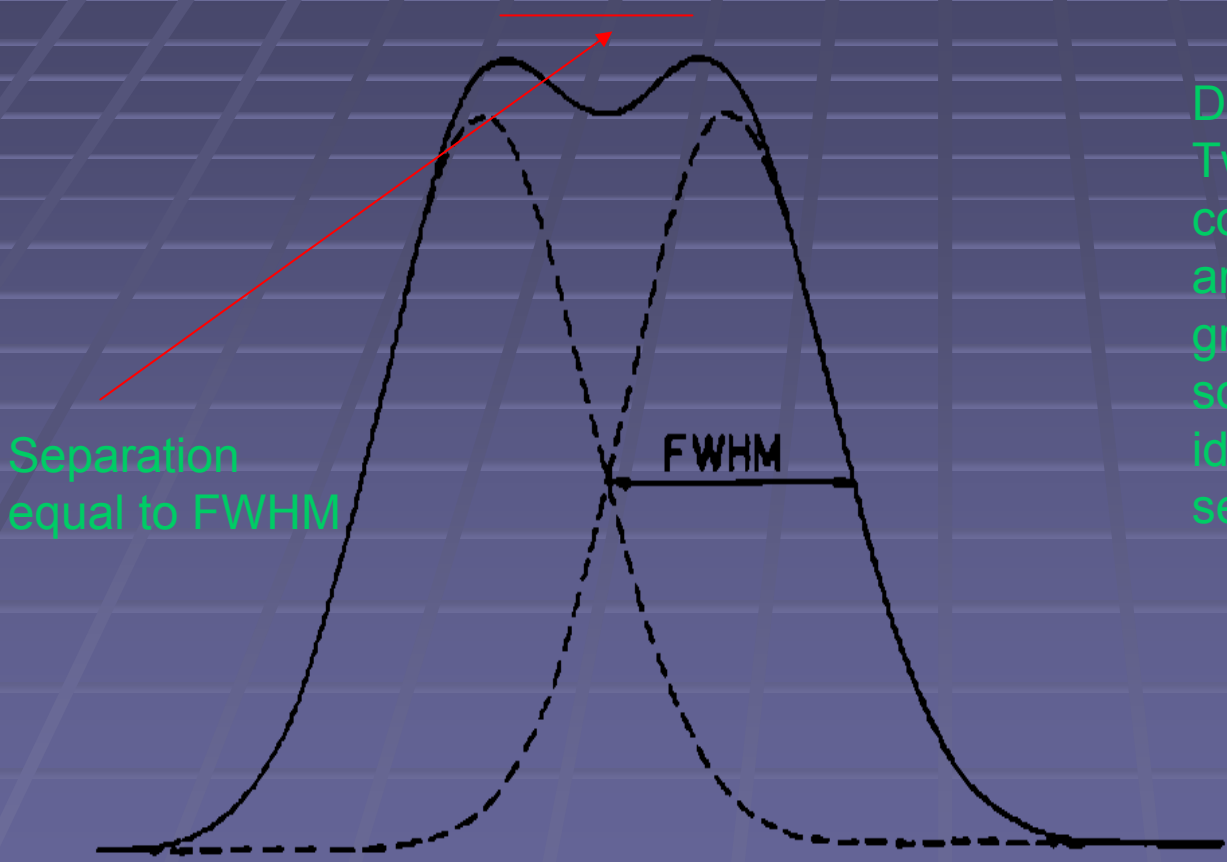
- **Cross-section** and detector **mass** → determine **probability** of energy conversion
- Charged particles → highly ionizing
- Neutral particles → indirect ionization
- The detector **mass** → depends on **radiation type**
- **Lower limit** is determined → by **noise**
- **Material** covering the entering **window**

Detector response

- Information of **energy** \longrightarrow the **amount** of ionization is proportional to **energy loss**
- Large detector volume completely absorb radiation \longrightarrow ionization gives energy information
- Information preserved or not, depends on design
- The amount of ionization \longrightarrow form of signal
- Signal: The **integral of pulse** with respect to **time**
- **Relation** between radiation **energy** and **pulse height** \longrightarrow **response** of detector
- Ideally this relation \longrightarrow linear
- A detector is **linear** for one radiation **but not** for other

Energy resolution

- For design consideration \longrightarrow energy resolution **most** \longrightarrow **important** factor
- Extent \longrightarrow to which a detector can **distinguish two** \longrightarrow close lying **energies**
- Resolution is measured \longrightarrow by sending a beam of **monoenergetic** particles
- Ideally a sharp delta function peak is expected
- Practically \longrightarrow a peak with **finite width** is observed, usually gaussian
- Width arises \longrightarrow because of **fluctuations** in the **number** of ionization and excitation



Defenition of energy resolution. Two peaks are generally considered to be resolved if they are separated by a distance greater than their FWHM. The solid line shows the sum of two identical Gaussian peaks separated by just this amount

- The resolution usually \longrightarrow in Full Width at Half Maxima (**FWHM**)
- Energies **closer** than this interval \longrightarrow normally **unresolvable**

$$\text{Resolution} = \frac{\Delta E}{E}$$

- **Nal** has **9%** resolution for **gamma** particles of about **1 MeV**
- **Germanium** detectors have **0.1%**
- Generally \longrightarrow resolution is **function** of energy, with \longrightarrow **ratio** improving with higher energy

- Due to \longrightarrow Poisson or Poisson like statistics of ionization and excitation
- Average energy required to produce an ionization is a fixed number, w
- For deposited energy, $E \longrightarrow$ Average ionization $J = E/w$
- Energy increase \longrightarrow ionization events increase \longrightarrow smaller relative fluctuations
- To calculate fluctuations \longrightarrow two cases

- Two cases
- Radiation energy is \longrightarrow not totally absorbed
- Thin transmission detector $\longrightarrow \frac{dE}{dx} \longrightarrow$ number of signal producing reactions \longrightarrow Poisson dist.

variance $\longrightarrow \sigma^2 = J$ $R = 2.35 \sqrt{w/E}$

- Full energy of radiation \longrightarrow absorbed
- Ionization events are not independent
- Poisson statistics is \longrightarrow not applicable

$\sigma^2 = FJ$ $R = 2.35 \sqrt{Fw/E}$

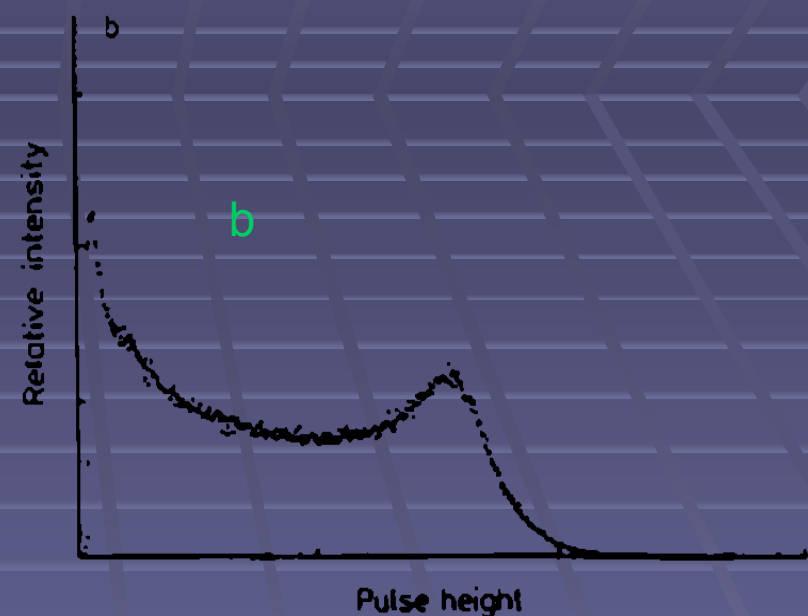
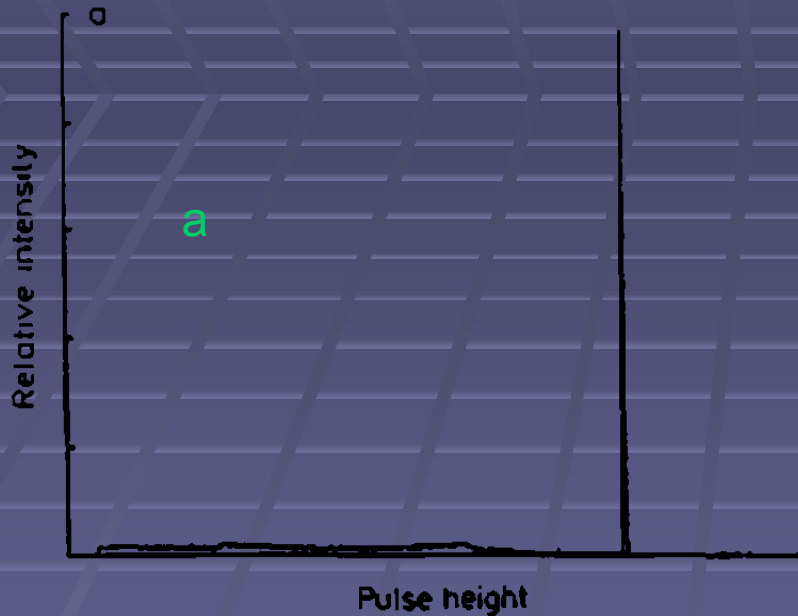
$F = \text{Fano Factor}$

The response function

- For the measurement of **energy spectra** , **response function** is crucial parameter
- Spectrum of pulse heights when bombarded by monoenergetic beam
- **Dirac delta function**
- Gaussian peak is not always realized, specially in case of neutral particles
- The **response function** is determined by **interactions** and design and **geometry** of detector

- Beam of monoenergetic electrons, incident on a detector, thick enough to stop particles.
- Some electrons will scatter out of detector
- **Low energy tail**
- Some **electrons** emit bremsstrahlung photons, which may escape from detector
- Events at a **lower energy than peak**
- Response function is **Gaussian peak** with low **energy tail** determined by amount of **scattering** and **bremsstrahlung**
- Response function can be **improved** by **changing design and geometry**

- A material of **small Z**, to **minimize backscattering and bremsstrahlung**
- **Relative intensity** of each structure in spectrum is determined by **relative cross-section** for each interaction mechanism
- **Response function** will be **different** at different energies for **different detector media**



- The response function of two different detectors for 661 keV gamma rays.
- a-) germanium detector which has large photoelectric effect cross-section relative to Compton scattering cross-section at this energy. A large photo peak with a small continuous Compton distribution.
- b-) response of organic scintillator with low Z. Compton scattering is predominant and only this distribution is visible

Response time

- The **time** taken by detector to **form** the **signal** after the arrival of particle
- **Duration** of signal is very important
- During this time **second event cannot be entertained**
- This contribute to dead time
- Limits the count rate

Detector efficiency

- Absolute efficiency

$$\mathcal{E}_{tot} = \frac{\text{events registered}}{\text{events emitted by source}}$$

- Function of detector geometry and probability of interaction

$$\mathcal{E}_{tot} = \mathcal{E}_{int} \mathcal{E}_{geom}$$

- Intrinsic efficiency

$$\mathcal{E}_{int} = \frac{\text{events registered}}{\text{events impinging on detector}}$$

- Intrinsic efficiency is function of radiation type, energy and detector material
- For charged particles generally intrinsic efficiency is good
- Problem of efficiency is important for neutral particles
- ϵ_{geom} is fraction source radiation which is geometrically intercepted by detector
- Angular distribution of incident radiation

Dead time

- Finite time to process the event
- Usually related to duration of pulse height
- Detector may or may not sensitive
- Insensitive detector → events lost
- Sensitive → events pileup
- Distortion and subsequent loss of information from both events
- Affect observed count rate
- Distort time distribution between arrival of events
- To avoid large dead time, rate should be small
- **RPC, rate capability 1 kHz/cm², dead time \approx 1 ns**
- Probability of another event is small

Introduction

- Gaseous detectors
- Greater mobility of electrons
- Obvious medium
- Charged particles detection
- Particle information easily transformed
- Efficiency
- Time resolution
- Rate capability
- cheap

Ionization mechanisms

- Energy loss → excitation, ionization



- Excitation, → resonant reaction, exact amount of energy transfer
- Cross-section → $\delta = 10^{-17} \text{ cm}^2$
- No free electron-ion pair



- Primary ionization, secondary ionization
- Penning effect → metastable states spin-parity difference → unable to deexcite → immediately
- Molecular ion → interacts neutral atom

- Ionization reactions → statistical in nature
- Two particles → not same → electron-ion pair
- For gases 1 e-i pair → 30 eV
- Average energy determines efficiency and time resolution of detector
- Recombination and electron attachment
- Electron affinity

$$R = 2.35 \sqrt{\frac{F_w}{E}}$$

Energy resolution of a particle

Working principle

- Ionization mechanism
- Working under different voltage conditions
- Exploiting different phenomenon
- Gas cylinder
- Noble gas
- Mean number of electron-ion pair created
- Recombination and electron attachment

Simple gaseous detector

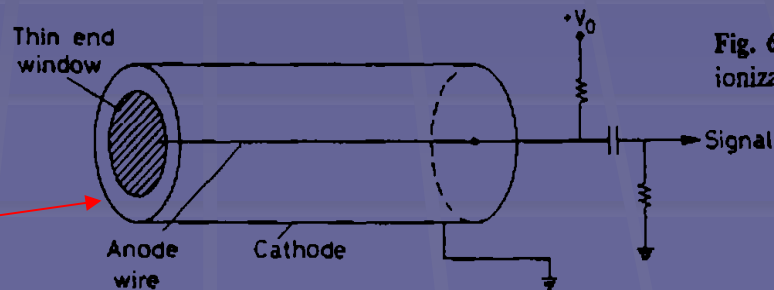


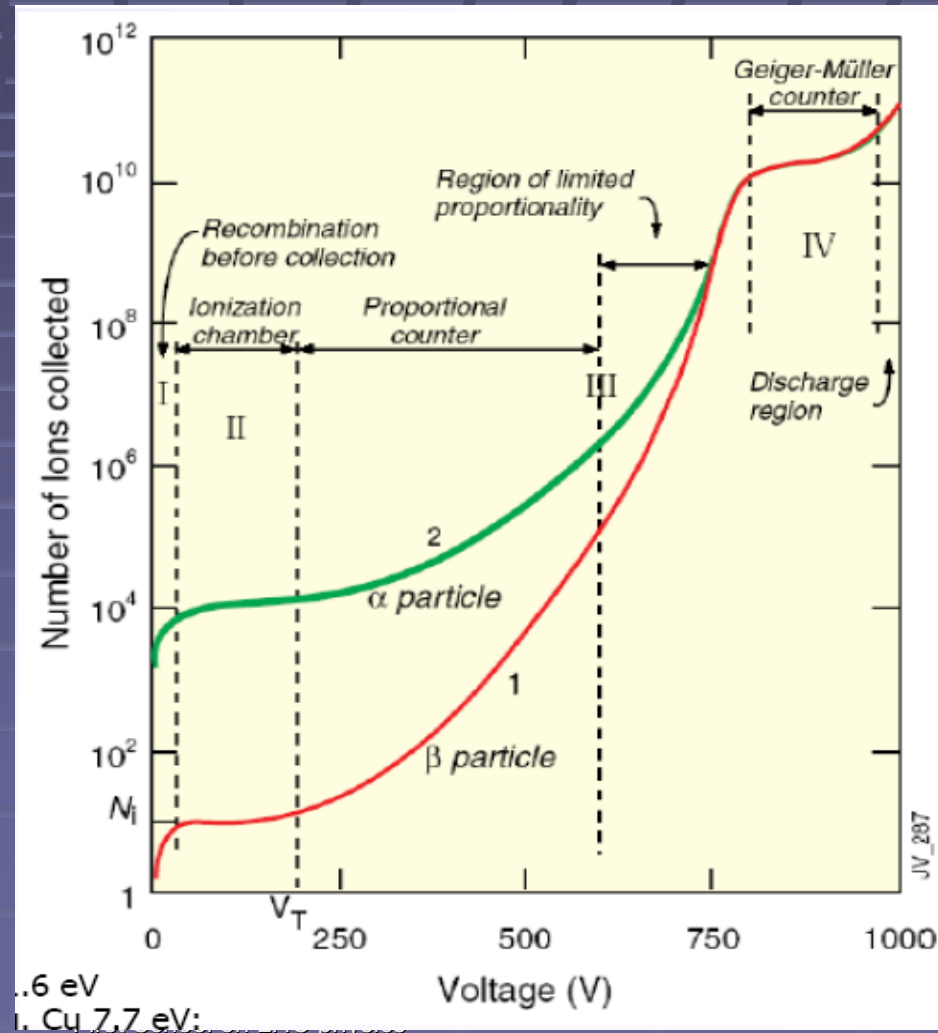
Fig. 6.1. Basic construction of a simple ionization detector

- Radial electric field $E = \frac{1}{r} \frac{V_0}{\ln\left(\frac{b}{a}\right)}$
- Electron-ion pair produced
- Direct ionization
- Indirect ionization
- Mean number of pair produced directly proportional to energy deposited by particle
- At zero voltage, recombination
- Recombination forces overcome, current increase
- All created pair collected, voltage increase no effect

- Detector working in this range → ionization chamber
- Voltage increase beyond region II
- Strong electric field → accelerate freed electrons → ionization
- Electrons liberated in secondary ionization accelerate → more ionization
- Ionization **avalanche** or **cascade**
- Strongest electric field near anode wire
- Number of electron-ion pair directly proportional to primary electrons

- Proportional amplification → as high 10^6
- Working in this range → Proportional chamber
- Further voltage increase
- Space charge distorts electric field about anode
- Proportionality begins to lost
- Region of limited proportionality
- Further V increase → discharge
- Chain reaction of many avalanche
- Secondary avalanche caused by → photons emitted by deexciting molecules

- Output current completely saturated
- Same amplitude regardless of particle energy
- To stop discharge → quench gas
- Working is this range → Geiger-Muller or break-down
- Further voltage increase → continuous breakdown



The total charge collected as function of V

0.6 eV
Cu 7.7 eV:

Conclusions

- Interaction of photons with material
- Detector parameters
- **Sensitivity**
- Detector response
- **Energy resolution**: The Fano Factor
- The response **function**
- **Response** time
- Detector **efficiency**
- **Dead** time
- Tried to understand basic **particle-material interactions**
- Working **principal** of **gaseous** detectors